

多重度分布とQCD相転移

中村純

nakamura@an-pan.org

共同研究者

永田桂太郎、森田健司

松本

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昨日の資料（コ一

- <http://home.riise.hiroshima-u.ac.jp/~nakamura/panflute/LTKf90v3.tar.gz>

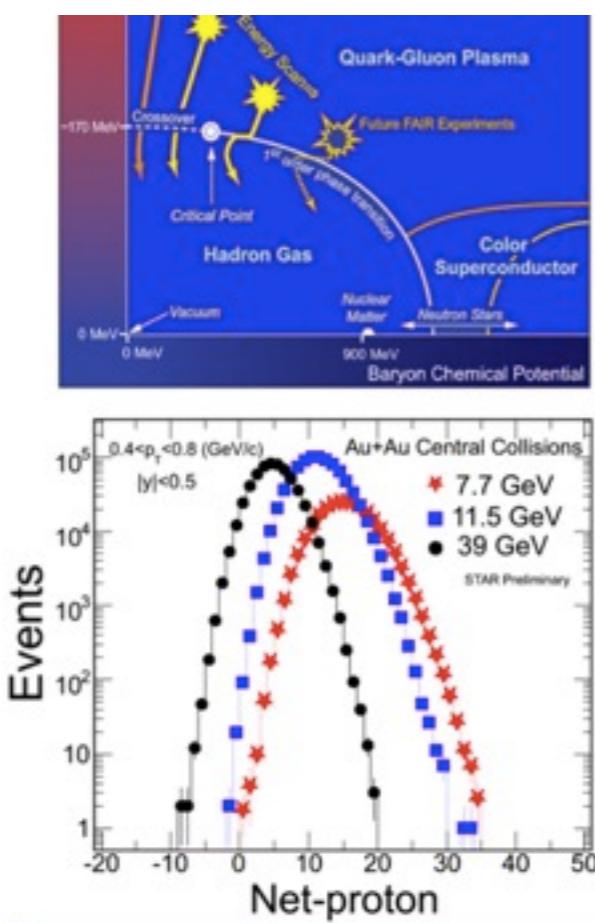
- そもそもは、高エネルギーの現象論をやっていて（博士論文は高エネルギー ハドロン原子核反応）、その道具として格子を始めました
- それからxx年

XQCD2012(2012 Aug. ワシントン)で

Nu Xuさんの話を聞きました

QCD Structure I (2012 Oct. 武漢)で

Multiplicityの話をたくさん聞きました



Nu Xu

"Extreme QCD", The George Washington University

- $\langle (\delta N)^2 \rangle \approx \xi^2$, $\langle (\delta N) \rangle \approx \xi$
- 3) Direct comparison:
 $S * \sigma \approx \frac{\chi_B^3}{\chi_B^2}$,
- 4) Extract susceptibility, temperature. A thermal equilibrium

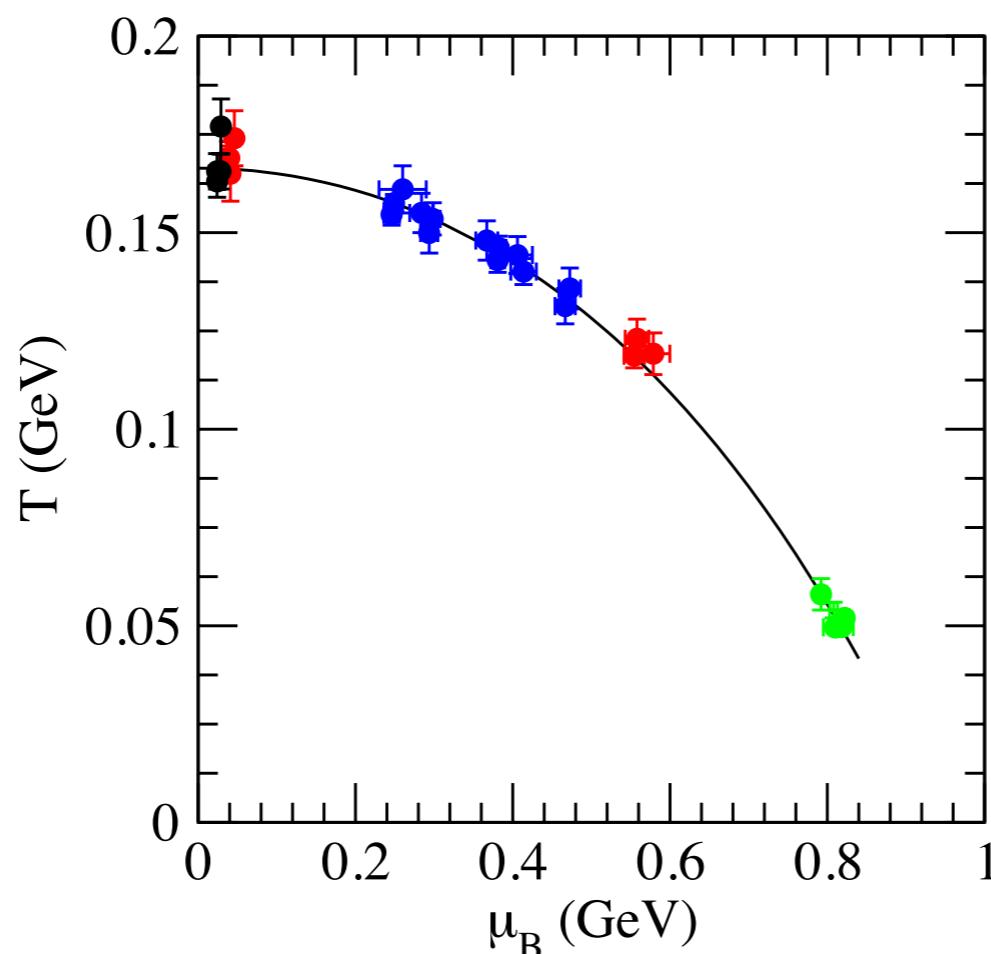
- A. Bazavov et al, PRD
- STAR Experiment: [PRD](#)
- M. Stephanov: [PRD](#)
- R.V. Gavai and S. Gupta: [PRD](#)
- S. Gupta, et al., [Science](#)
- F. Karsch et al, [PLB](#)
- M. Cheng et al, [PRC](#)
- Y. Hatta, et al, [PRL](#)

あつ、
Multiplicity
だ！



Freeze-out

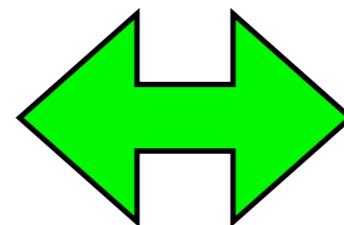
Temperature and Chemical Potential
when the particles were born
were estimated.



Cleymans et al.

Phys. Rev. C73:034905, 2006

Canonical Partition Function



Grand-Canonical Partition Function

$$Z(\xi, T) = \text{Tr } e^{-\beta(H - \mu \hat{N})} \quad (\beta = 1/T)$$

$$= \sum_n \langle n | e^{-\beta(H - \mu \hat{N})} | n \rangle$$

$$= \sum_n \langle n | e^{-\beta H} | n \rangle (e^{\mu n / T})$$

If $[H, \hat{N}] = 0$

$$= \sum_n \langle n | e^{-\beta H} | n \rangle (e^{\mu n/T})$$

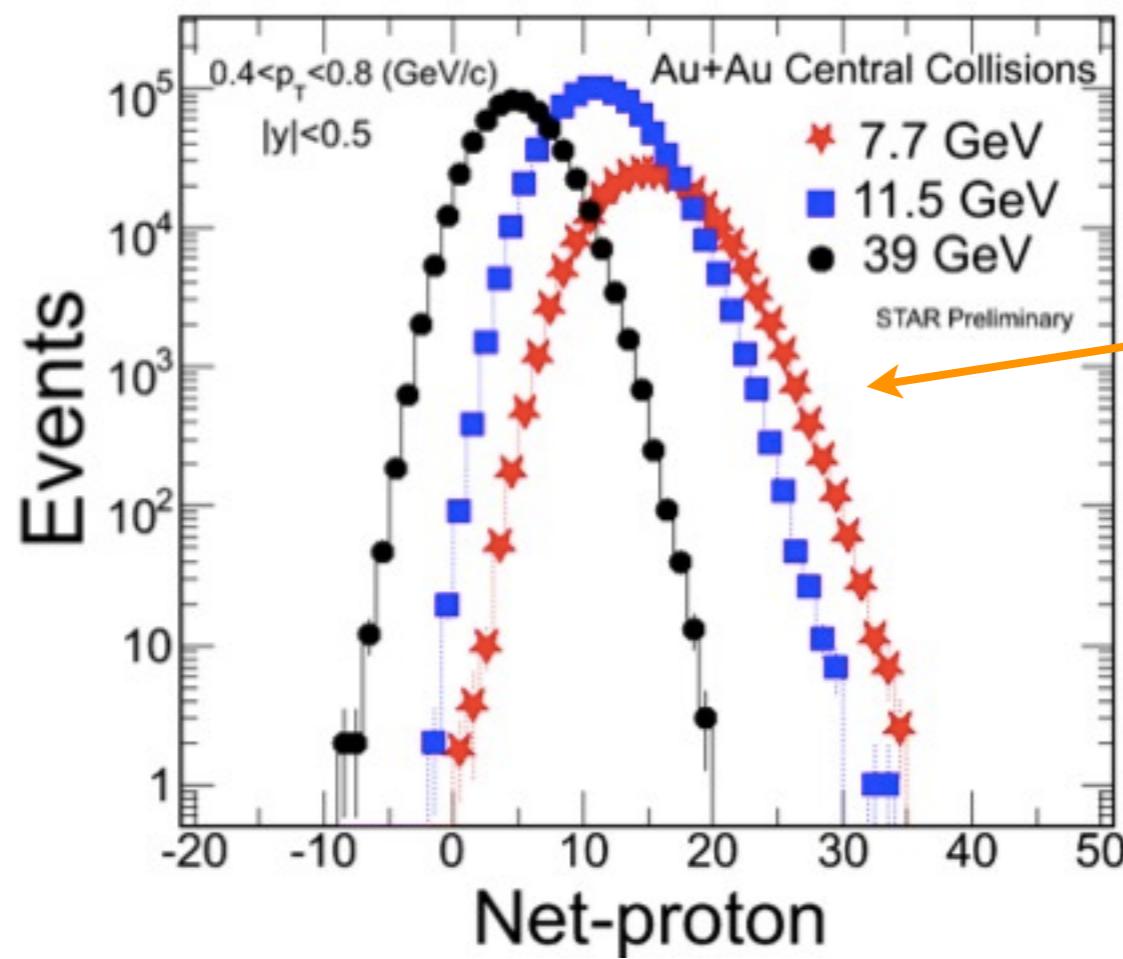
$$Z(\xi, T) = \sum_n Z_n(T) \xi^n$$

where $Z_n = \langle n | e^{-\beta H} | n \rangle$
 $\xi = e^{\mu/T}$ (Fugacity)

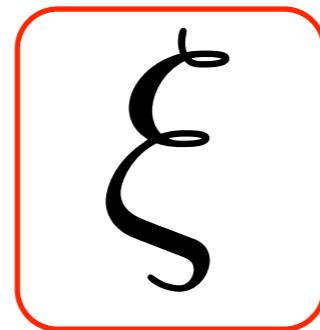
$$Z(\xi, T) = \sum_n Z_n(T) \xi^n$$



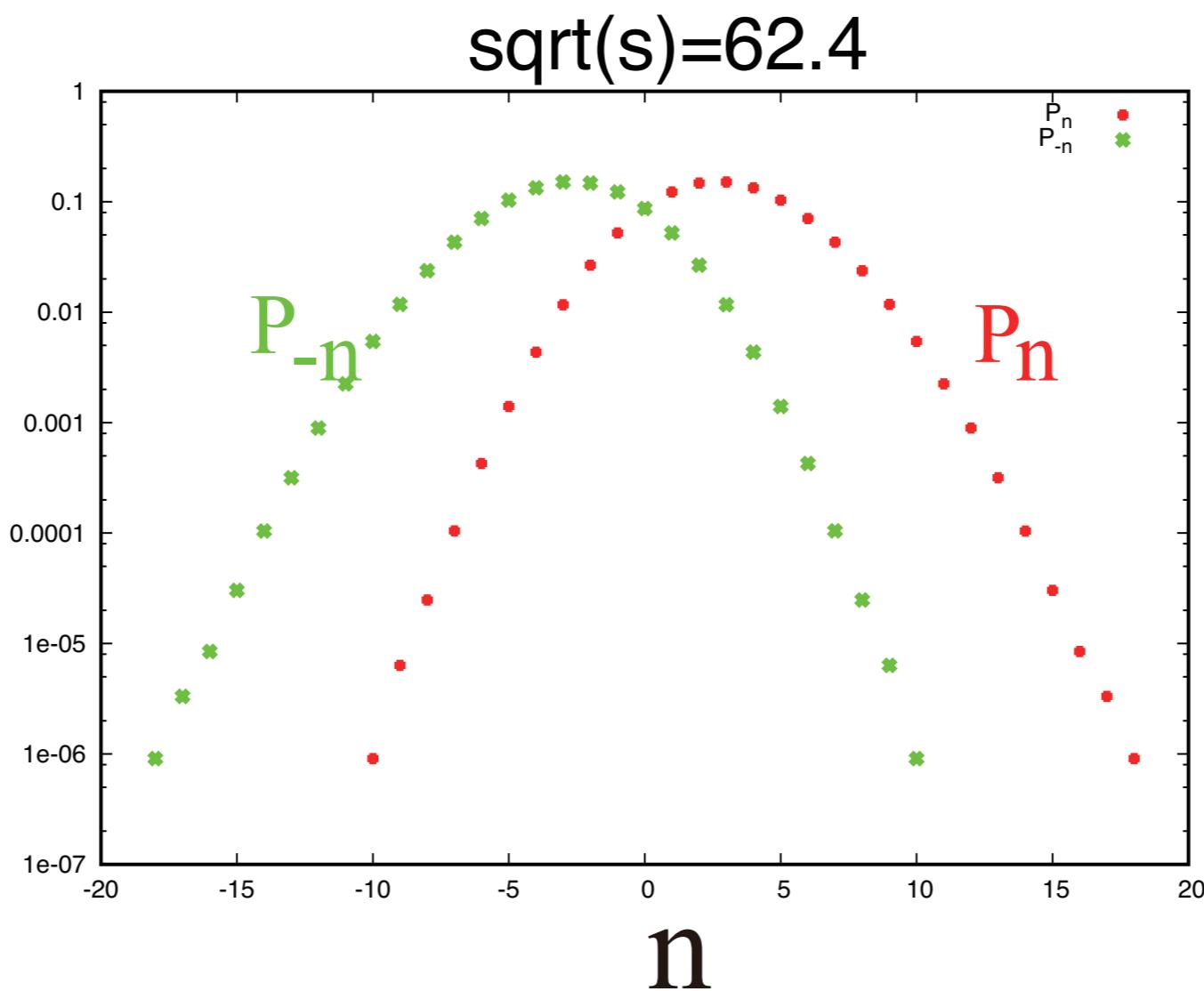
Partition Function is
Sum of the Probability ...
If I know ξ , then I have Z_n .



How can I obtain Zn from $P_n = Z_n \xi^n$?



unknown ?!



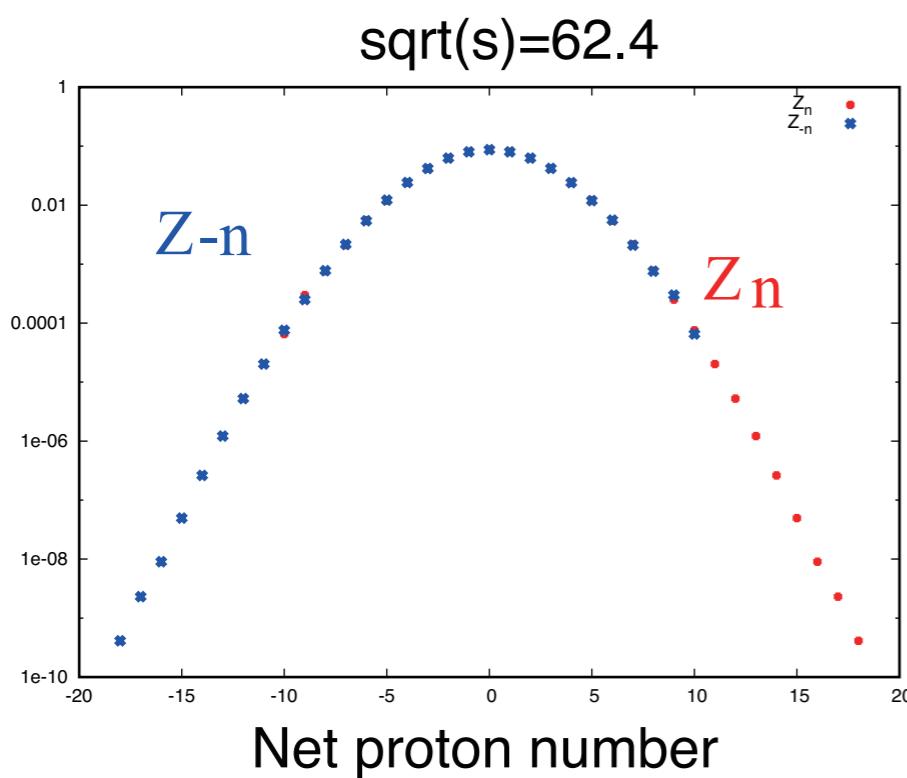
$$Z_n = P_n / \xi^n$$

Particle-AntiParticle
Symmetry !

$$Z_{+n} = Z_{-n}$$

$$\xi = e^{\mu/T}$$

One Parameter !

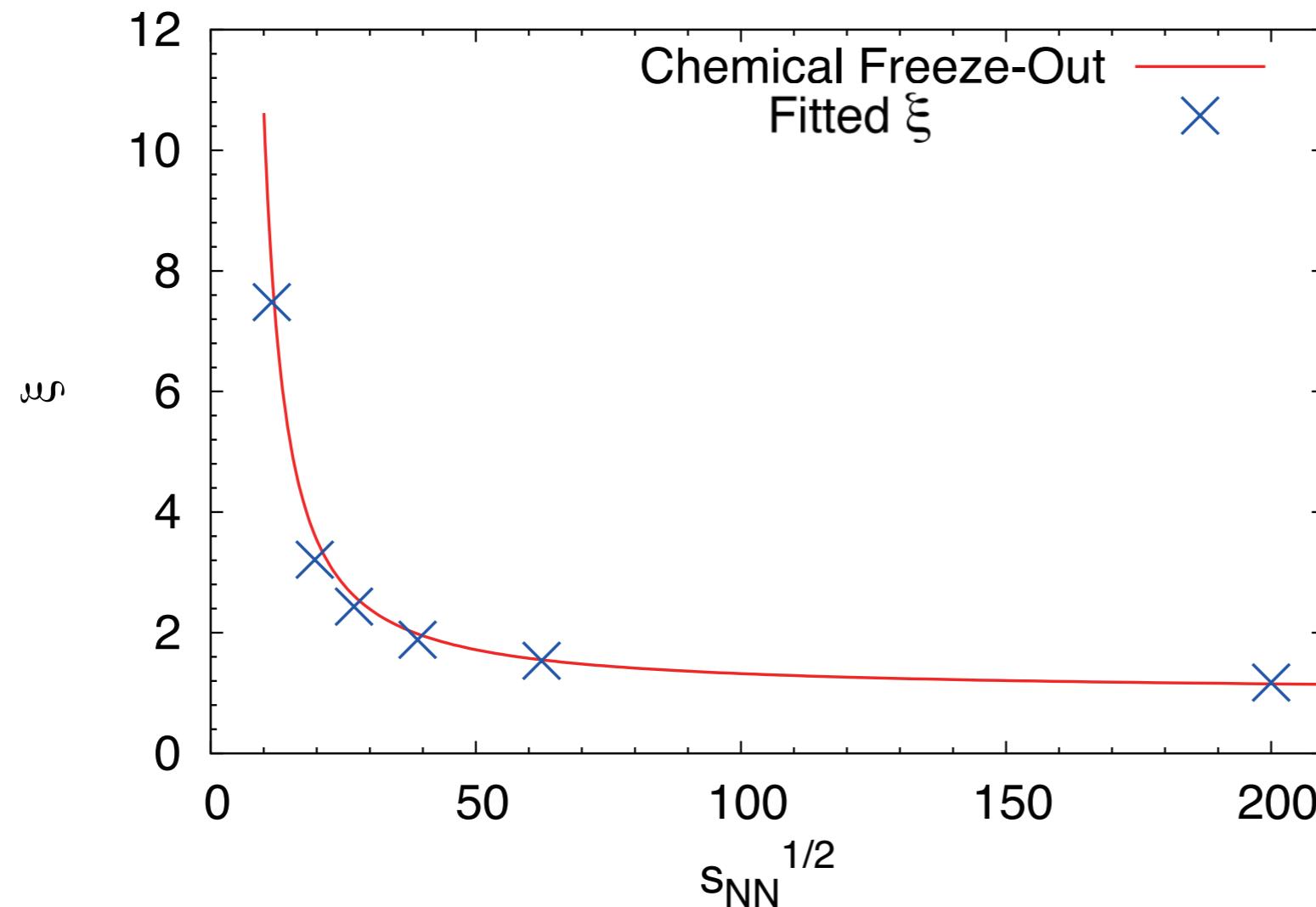


Fit ξ

Potato -UntiPotato
Symmetry ?



Fitted ξ are very consistent with those by Freeze-out Analysis.



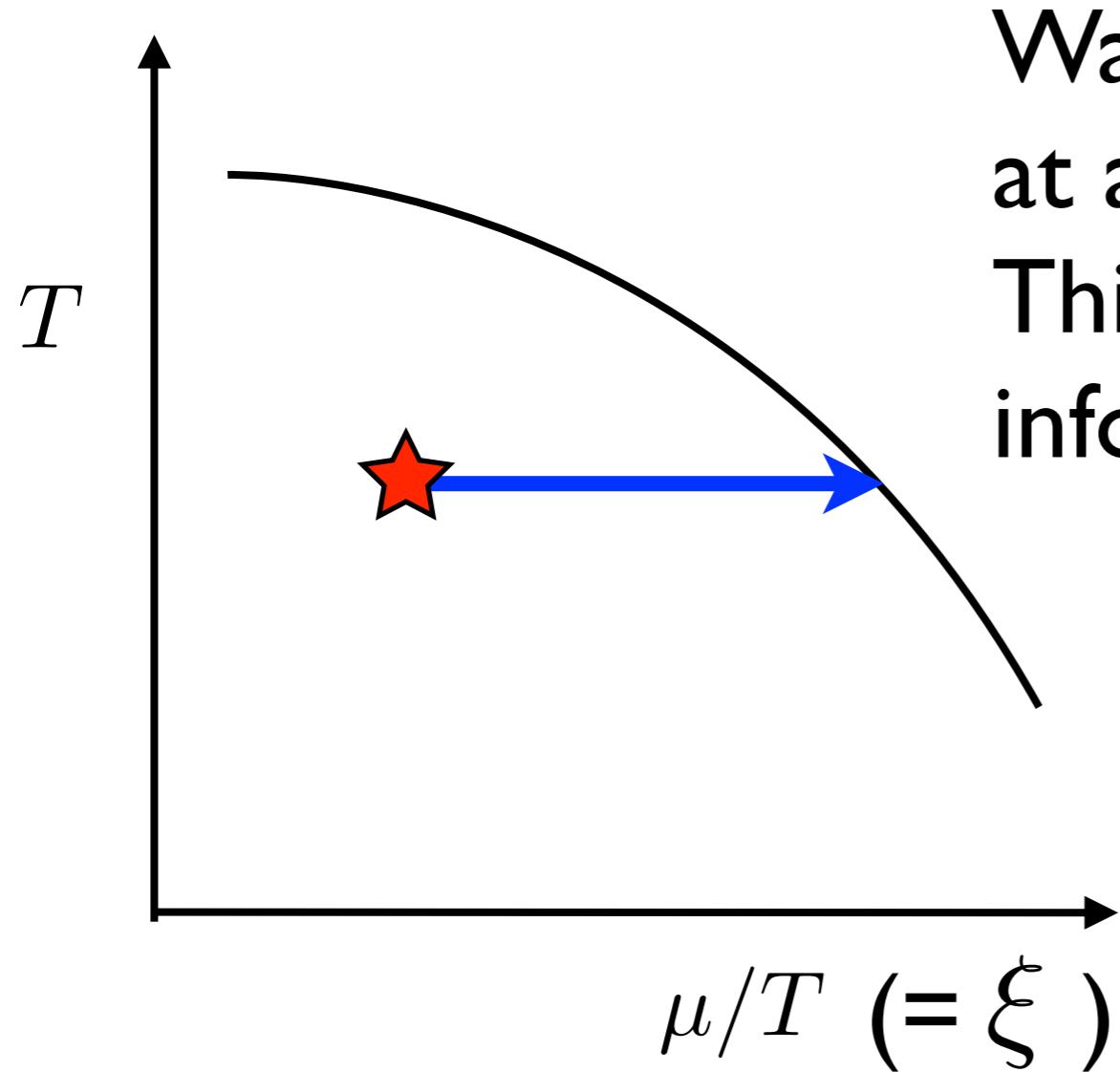
\times This work

— Freeze-out

J.Cleymans,
H.Oeschler,
K.Redlich and
S.Wheaton
Phys. Rev. C73,
034905 (2006)

$$Z(\xi, T) = \sum_n Z_n(T) \xi^n$$

Now we have Z_n of RHIC data
($\sqrt{s}=9.6, 27, 39, 62.4, 200$ GeV)

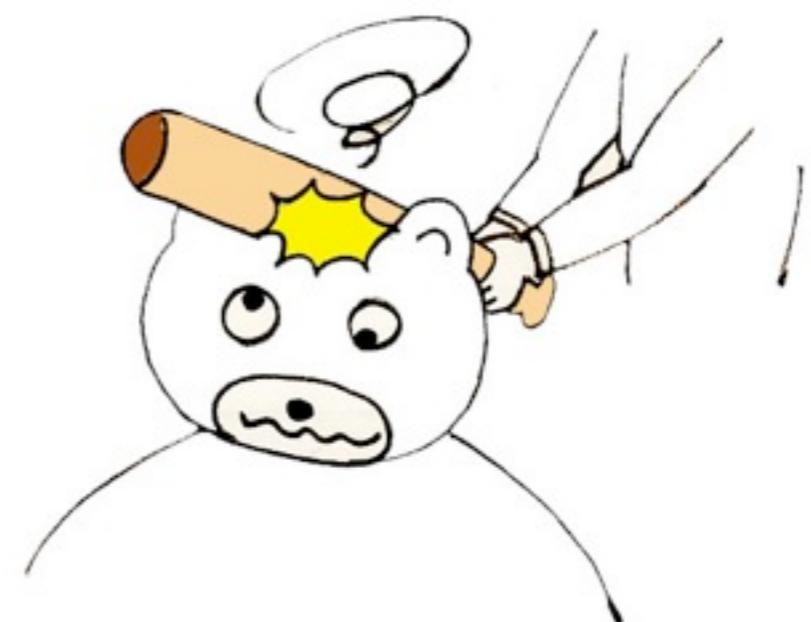


Wao ! We can calculate
at any density !
This includes all QCD Phase
information !



$$Z(\xi, T) = \sum_{n=-N_{max}}^{+N_{max}} Z_n(T) \xi^n$$

Do not forget that your n is finite !

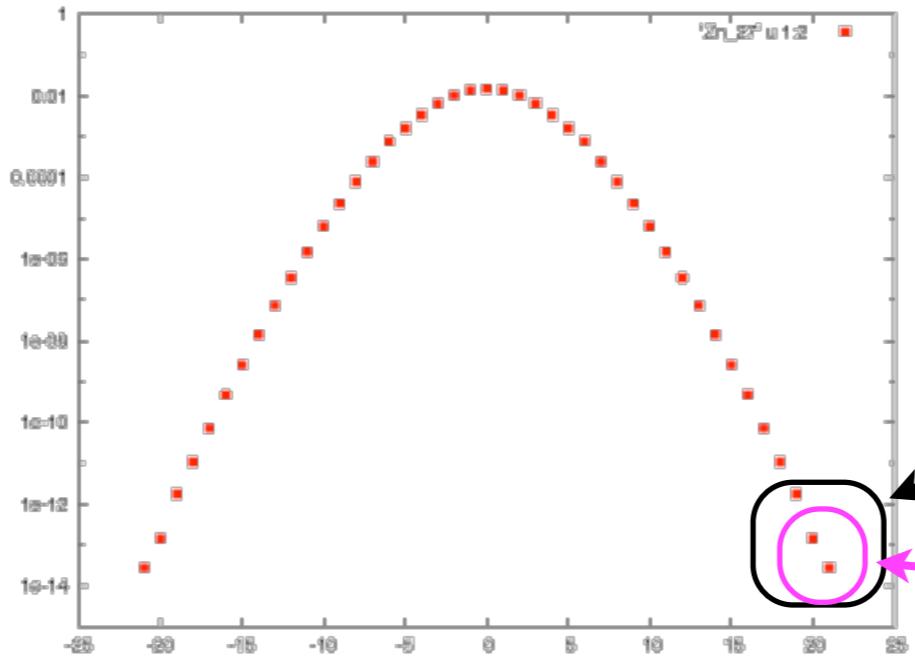


Hunting the QCD Phase Transition Regions

Find Rooms where No Criminal.
→ The Target is in other Room.

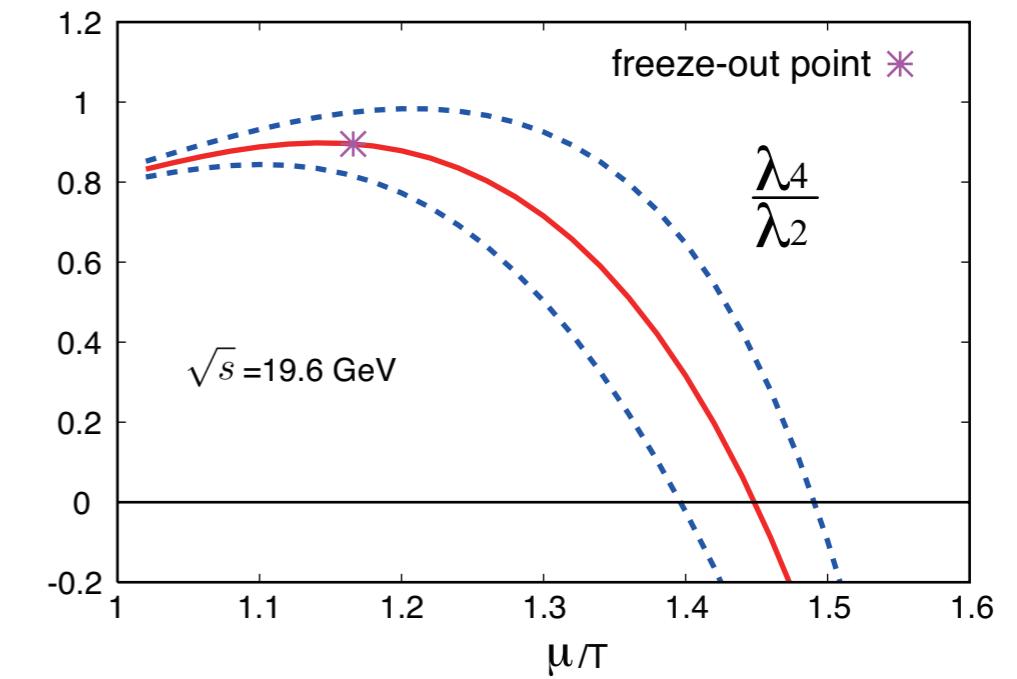
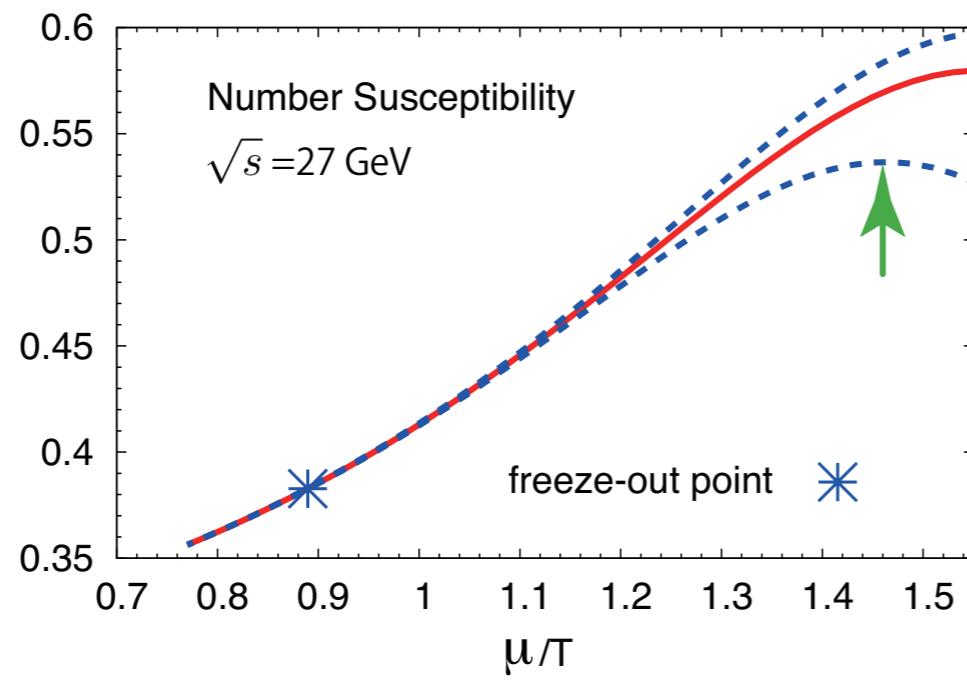


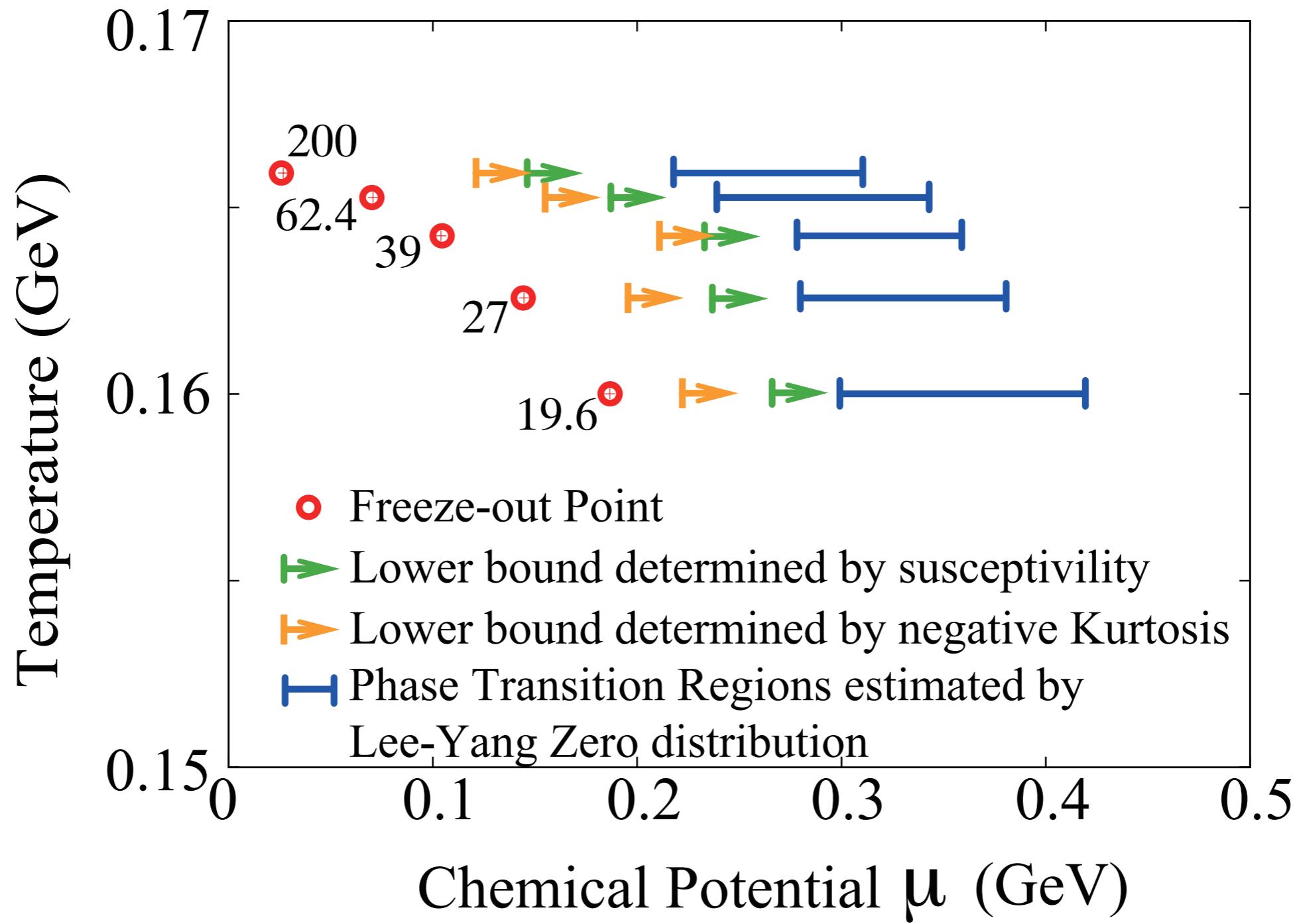
Not here ! Then, ..



What happens
if we increase
these points 15%

if we drop
these points





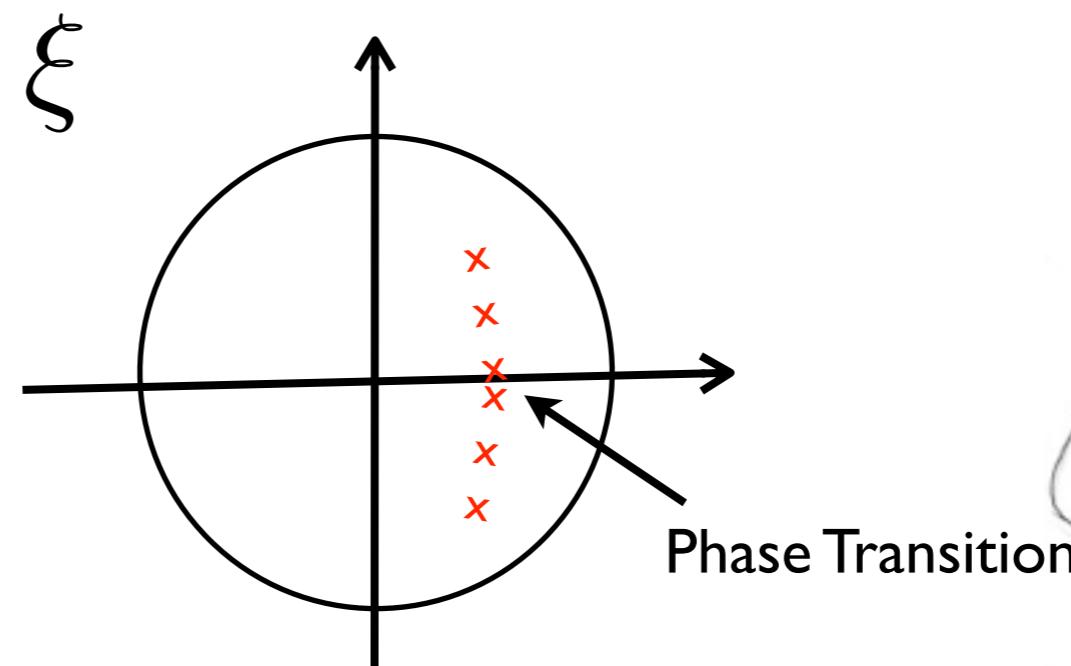
Lee-Yang Zeros

Zeros of Z_n in Complex Fugacity Plane.

$$Z(\alpha_k) = 0$$



Great Idea to investigate
a Statistical System



Simple Factorization Problem ?

$$Z(\xi, T) = \sum_{n=-N_{max}}^{+N_{max}} Z_n(T) \xi^n = \prod_k (\xi - \alpha_k) = 0$$

For large Nmax, this is
an ill-posed problem.

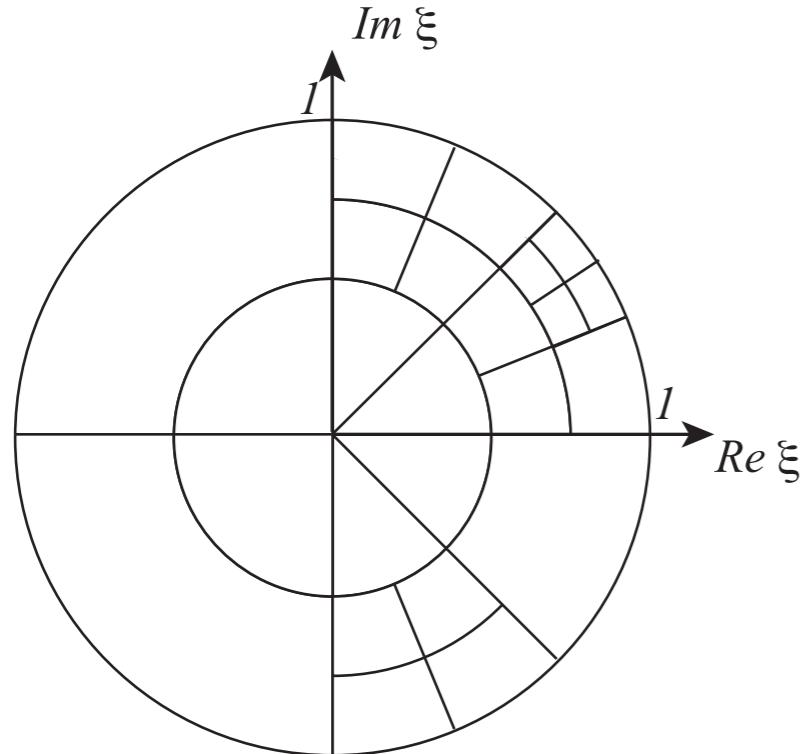


IMSL Library does not work.

Minimum Search of $|Z(\xi)|$

Difficult to confirm a real Zero from
very small value.

cut Baum-Kuchen (cBK) Algorithm



$$f(\xi) = \prod_k (\xi - \alpha_k)$$

$$\frac{f'}{f} = \sum_k \frac{1}{\xi - \alpha_k}$$

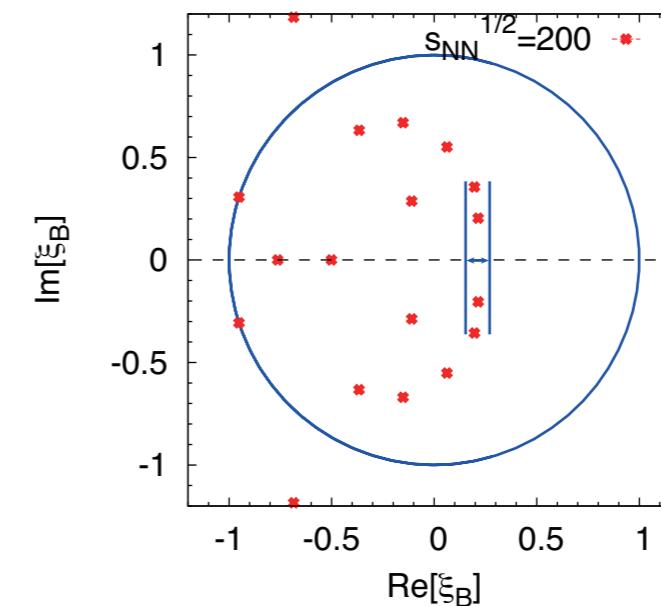
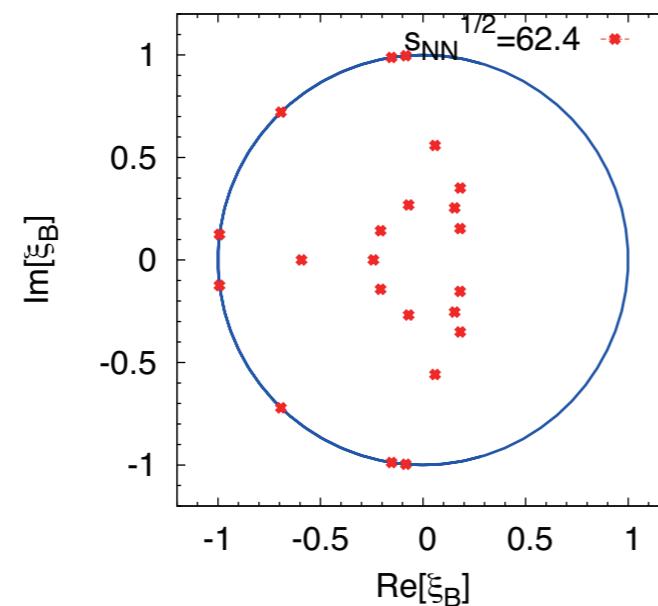
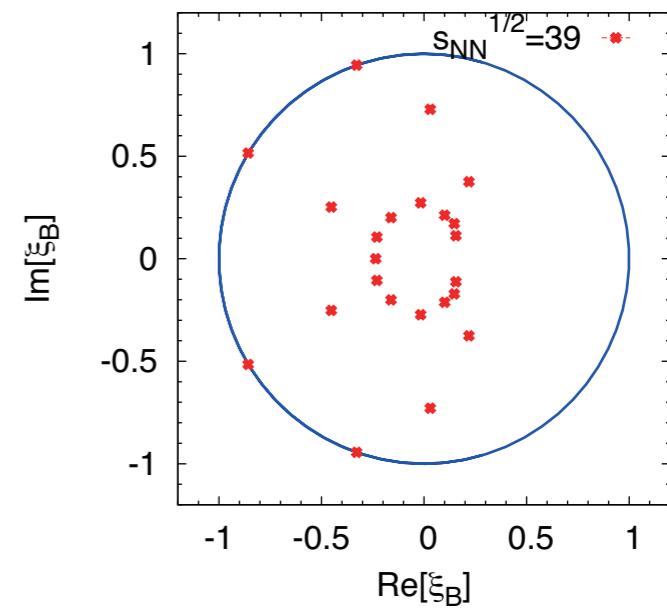
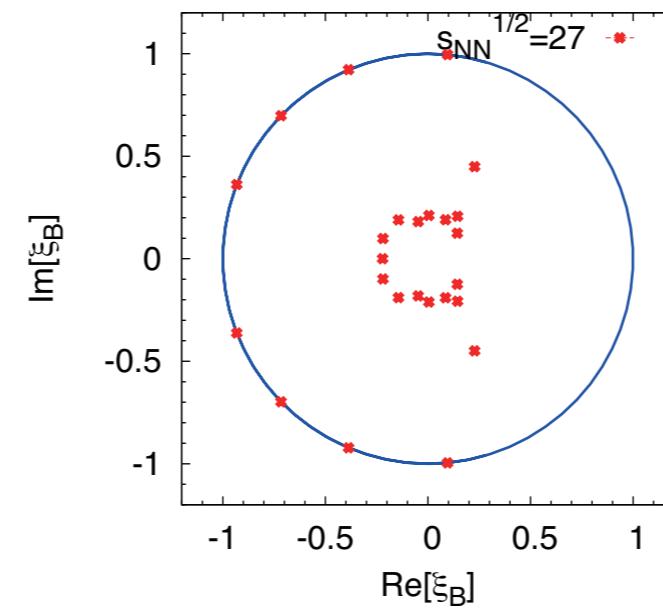
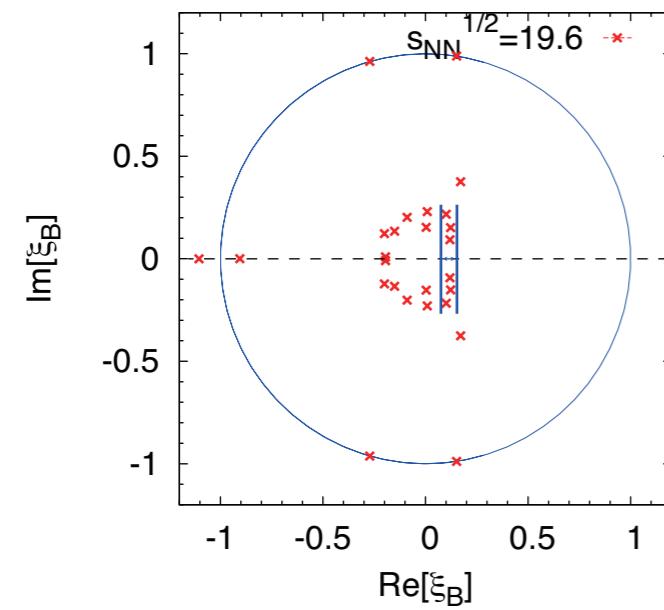
$$\frac{1}{2\pi} \oint_C \frac{f'}{f} d\xi = \text{Number of Zeros in Contour C}$$

50 - 100 number
of significant digits

A Coutour is cut into
four pieces
if there are zeros insied.



LYZ: RHIC Experiments



Lattice

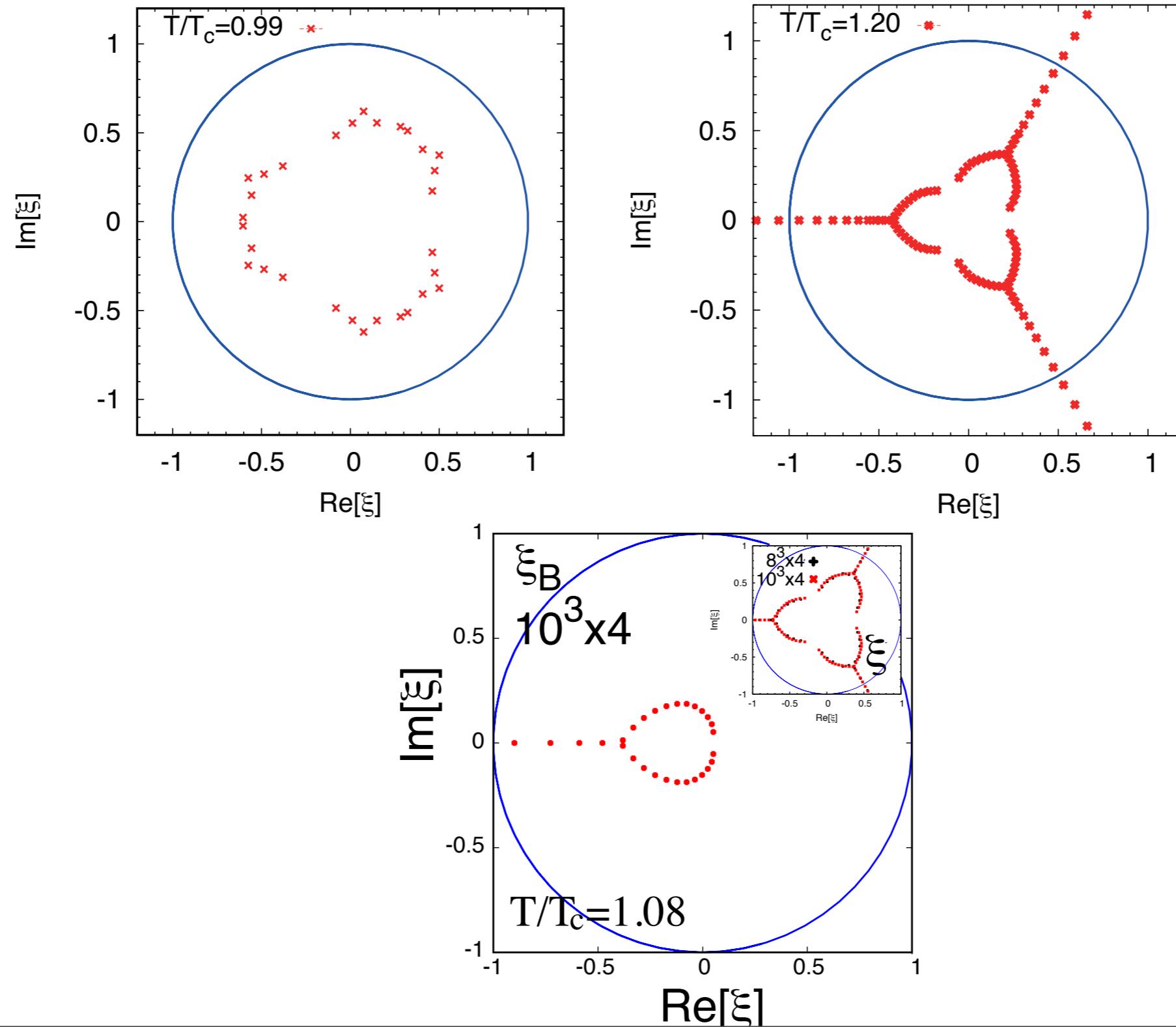
$$Z_{GC}(\mu) = \int DU \left[\frac{C_0 \sum c_n \xi^n}{\det \Delta(0)} \right]^{N_f} (\det \Delta(0))^{N_f} e^{-S_G}$$

$$\begin{aligned} & \int DU \left(\sum a_n \xi^n \right) (\det \Delta(0))^{N_f} e^{-S_G} \\ &= \int DU \left(\sum a_n \xi^n \right) (\det \Delta(0))^{N_f} e^{-S_G} \end{aligned}$$

$$= \sum_n \xi^n \int DU a_n (\det \Delta(0))^{N_f} e^{-S_G}$$

$$Z_{GC}(\mu) = \sum Z_n \xi^n$$

LYZ: Lattice QCD



Lattice: How to Calculate

$$Z(\mu, T) = \int \mathcal{D}U (\det \Delta(\mu))^{N_f} \exp(-S_G)$$

$$\det \Delta(\mu) \rightarrow \left(\frac{\det \Delta(\mu)}{\det \Delta(0)} \right) \det \Delta(0)$$

$$\det \Delta(\mu) \propto \xi^{-N_{\text{red}}/2} \prod_{n=1}^{N_{\text{red}}} (\lambda_n + \xi)$$

$$\propto \sum_{n=-N_{\text{red}}/2}^{N_{\text{red}}/2} c_n \xi^n$$

Fugacity Expansion
Nagata and A. Nakamura,
Phys. Rev. D82 094027

Skellam

Skellamの場合
はどうなるん
ですか？

大変だ、新しい話だ
勉強しないと！



Skellam

新しくないよ

Skellam, J. G. (1946)

"The frequency distribution of the difference between two Poisson variates belonging to different populations". *Journal of the Royal Statistical Society, Series A*, 109 (3), 296



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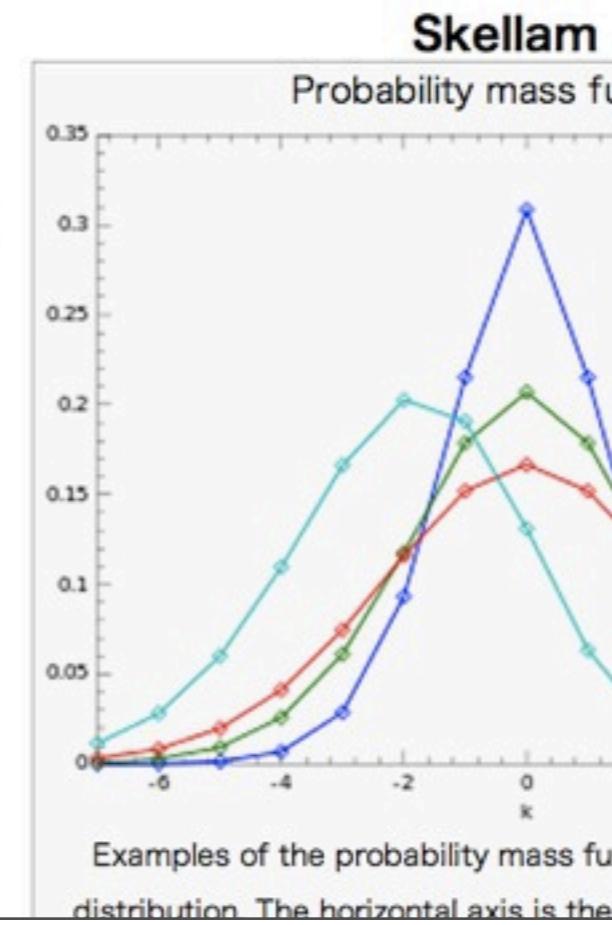
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Skellam distribution

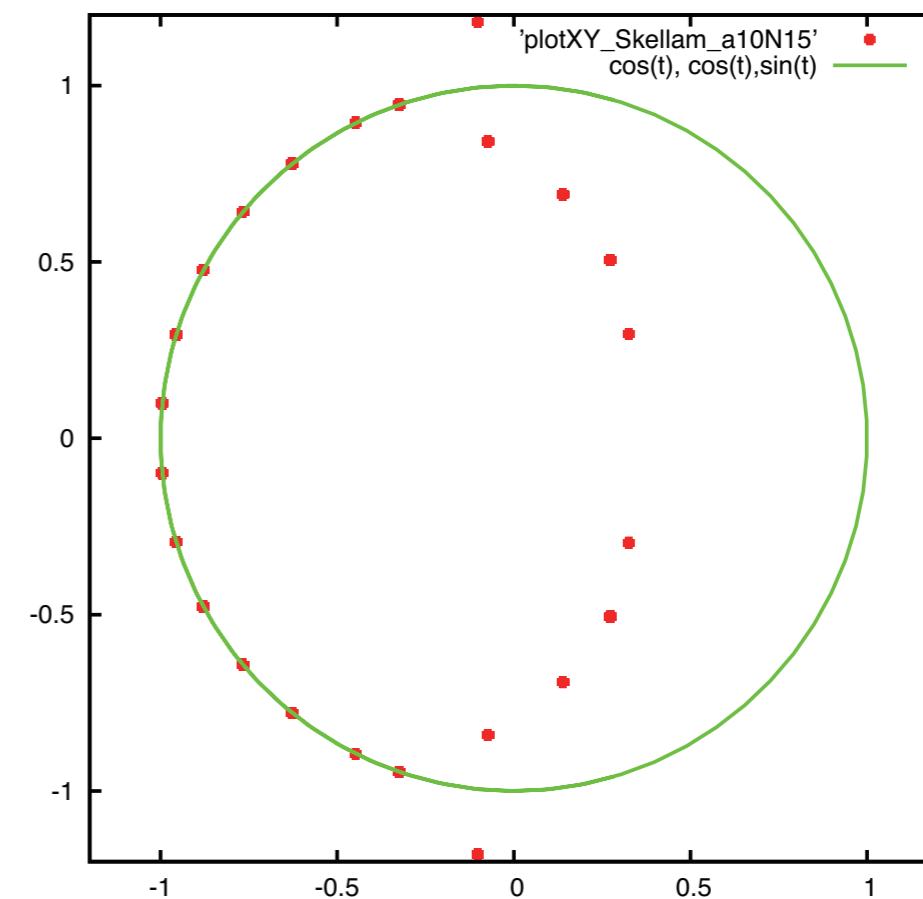
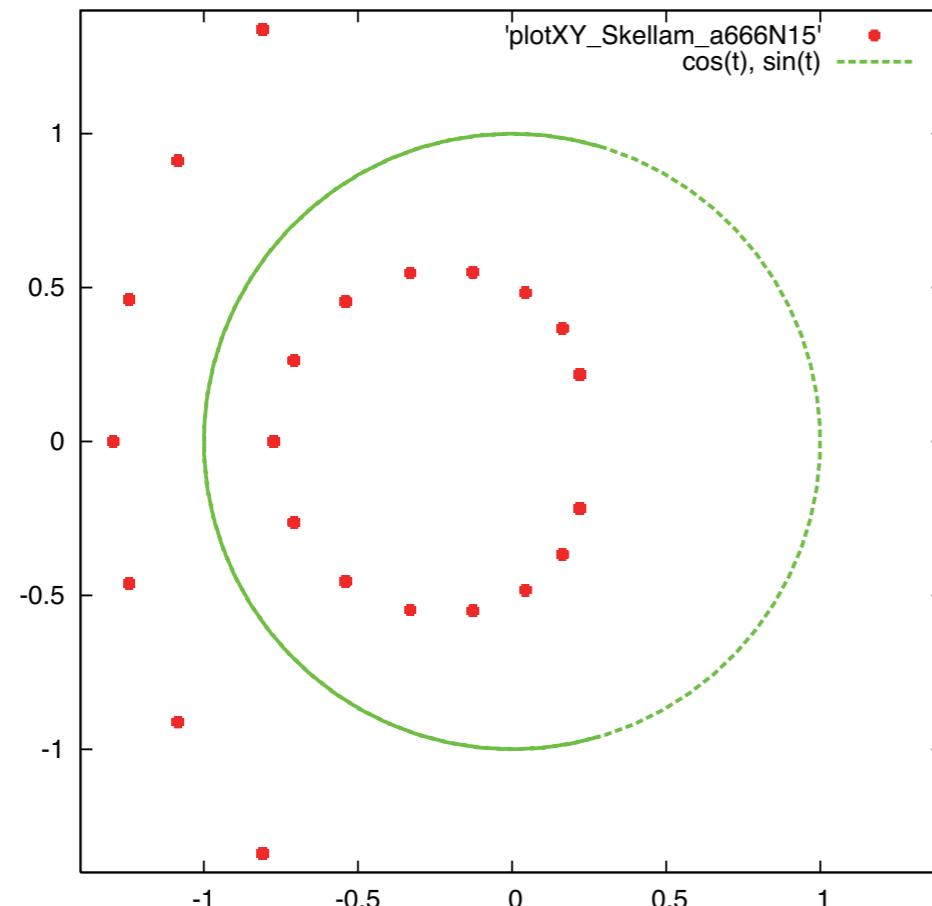
From Wikipedia, the free encyclopedia

The **Skellam distribution** is the discrete probability distribution of the difference $n_1 - n_2$ of two statistically independent random variables N_1 and N_2 each having Poisson distributions with different expected values μ_1 and μ_2 . It is useful in describing the statistics of the difference of two images with simple photon noise, as well as describing the point spread distribution in certain sports where all scored points are equal, such as baseball, hockey and soccer.

The distribution is also applicable to a special case of the difference of dependent Poisson random variables, but just the obvious case where the two variables have a common additive random contribution which is cancelled by the difference: see Karlis &



LYZ : Skellam



Conclusion

- メチャクチャ楽しい仕事だった
- 突っ込みどころはまだたくさん
 - proton multiplicityは保存量じゃないだろう
 - Lee-Yang zeros分布とO(4)スケーリングの関係は？
 - Latticeは $T < T_c$ がちゃんと計算できていない
- 実験でNmaxを大きくしてくれれば、QCD相図はかなり押さえ込まれる
- 格子で T_c 以下を計算しなければ
 - 虚数化学ポテンシャル領域の計算だが大きなノイズ